

Mechanics — Math Walk-Through

Purpose. Derive the working kernels for bending, travel-time delay, and clock drift in the VMS framework, and show the smooth-limit bridge to the standard weak-field results.

Conventions and Locks

Clarifications & Conventions (alignment with Narrative and Appendix)

- “Route action” is the display-area action $S[\text{path}] \equiv \int A_d ds$. We compare routes by S ; phase appears only when needed via $\varphi_{\text{phase}} = S/S_0$.
- Background/cost map $n(x)$: think of n as a smooth cost map tied to missing space; higher n means “this step hides a bit more display-area than average.” It is not a new substance.
- Link to potential and sign convention: set $n(x) = 1 + \eta_0 \varphi(x)$ with $\varphi = \Phi/c^2$ and $\Phi < 0$ near mass. With $|\eta_0|$ fixed by deflection ($|\eta_0| = 2$), bending is toward the mass and travel-time shifts are positive.
- Scope note: weak + strong field live here. For steep profiles $n(x)$ evaluate the same kernels piecewise; avoid turning points with refinement.
- Single dimensional scale: $S_0 = \hbar$ (SI-locked).
- Projected display-area: $A_d = \pi r^2$ (orthographic; no 4π shells).
- Route action: $S[\text{path}] = \int A_d ds$ (ratios first; phase = S/S_0 when needed).
- Speed of advance: c ; Newton’s constant: G .

Background Model and Ray Equation

A stationary mass is treated as a closed circulation whose far-field effect is summarized by a smooth cost map $n(x)$. Adopt the WKB/eikonal ansatz $U = A e^{i k_0 S}$. Collecting powers of k_0 yields the eikonal and transport equations:

$$|\nabla S|^2 = n(x)^2$$

$$2 \nabla S \cdot \nabla A + A \nabla^2 S = 0$$

Rays follow the normal to S . Fermat form $\delta \int n ds = 0$ gives the ray equation:

$$d/ds (n \mathbf{t}) = \nabla n \quad (\mathbf{t} \text{ is the unit tangent})$$

1. Bending Kernel

Take the component perpendicular to the ray:

$$d\theta/ds = (1/n) (\nabla_{\perp} n) \approx \nabla_{\perp} \ln n \quad (\text{small angles})$$

Integrate along the unperturbed straight line to get the net deflection vector:

$$\alpha^{\vec{}} \approx \int (\nabla_{\perp} \ln n) ds$$

Weak-field model: $n(\mathbf{x}) = 1 + \eta_0 \varphi(\mathbf{x})$, with $\varphi = \Phi/c^2$ and $\Phi = -GM/r$. For $|\eta_0 \varphi| \ll 1$, $\ln n \approx \eta_0 \varphi$, so:

$$\alpha^{\vec{}} \approx \eta_0 \int \nabla_{\perp} \varphi(\mathbf{x}) ds$$

For a grazing pass with impact parameter b :

$$|\alpha| \approx 2 |\eta_0| GM / (b c^2)$$

Matching the standard weak-field value $|\alpha| = 4 GM/(b c^2)$ fixes $|\eta_0| = 2$ (sign chosen to match deflection toward mass).

1A. Strong-Field: Photon-Sphere and Capture (dictionary check)

For steep, spherically symmetric profiles $n(r)$ that grow sharply as r decreases, closed circular null routes appear when transverse curvature balances forward advance. Define the critical impact parameter b_c as the smallest value for which a passing route still escapes; as $b \rightarrow b_c^+$, the deflection $|\alpha|$ diverges and multiple “looped” images arise. For $b < b_c$, routes are captured.

In the GR-limit dictionary this reproduces the familiar values (photon sphere and critical impact parameter):

$$b_c = 3\sqrt{3} GM / c^2$$

Evaluation uses the same kernel $\alpha^{\vec{}} = \int \nabla_{\perp} \ln n ds$ on the steep $n(r)$ profile; the divergence near b_c is a geometric property of the logarithmic integral in a rapidly varying cost map.

$$r_{ph} = 3 GM / c^2$$

2. Travel-Time Delay

Propagation time increment relative to empty space:

$$\Delta t = (1/c) \int [n(\mathbf{x}) - 1] ds \approx (\eta_0/c) \int \varphi(\mathbf{x}) ds$$

For a point mass on a straight-line path with closest approach b , evaluation gives the standard logarithmic form after using $|\eta_0| = 2$:

$$\Delta t \approx (2 GM / c^3) \ln((r_S + r_R + D) / (r_S + r_R - D))$$

2A. Observer Dictionary — Lensing Map and Time-Delay Decomposition

For far-field geometry (thin-lens approximation) with distances D_l (observer→lens), D_s (observer→source), D_{ls} (lens→source), connect the deflection kernel to observed angles:

Here $\alpha(\theta) \approx (1/D_l) \int \nabla_{\perp} \ln n ds$ written along the unperturbed path. Relative arrival time between two images i and j splits into geometric and potential parts:

The projected potential ψ is linked to $\ln n$ via the same delay kernel $\Delta t = (1/c) \int (n - 1) ds$. This dictionary is only a re-expression of the kernels above for inference use; no new parameters are introduced.

$$\Delta t_{ij} = (D_{\Delta} / c) [\frac{1}{2}(\theta_i - \beta)^2 - \frac{1}{2}(\theta_j - \beta)^2 - \psi(\theta_i) + \psi(\theta_j)]$$

$$\beta = \theta - (D_{ls}/D_s) \alpha(\theta)$$

r_S, r_R : source/receiver distances from the mass; D : separation.

3. Clock Drift (Redshift / Time-Dilation)

In a stationary background the local tick rate obeys, to leading order:

$$d\tau \approx dt \cdot \sqrt{1 + 2\phi} \quad \Rightarrow \quad \Delta v/v \approx -\Delta\phi$$

We treat this as the smooth-limit dictionary of the same cost picture (no new parameters).

4. Bound Orbits: Periapsis Precession (Sketch)

For slowly moving matter routes in a central profile $n(r) = 1 + \eta_0 \phi(r)$, transverse curvature from $\nabla_{\perp} \ln n$ produces a small advance of periapsis per orbit. In the weak-field, low-speed limit the result reduces to:

$$\Delta\omega \approx 6\pi GM / (a(1 - e^2)c^2) \quad \text{per orbit}$$

Full derivation is provided via an effective potential constructed from the ray equation (Appendix).

5. Rotators and Frame-Dragging (Preview)

With angular momentum J , the background acquires a small odd-parity correction (azimuthal bias). In the smooth-limit dictionary this corresponds to the Lense-Thirring rate:

$$\Omega_{LT} = 2GJ / (c^2 r^3) \quad (\text{axis-aligned, far field})$$

We treat these corrections as extensions; in bench contexts they are reported as bounds unless a positive detection is required.

6. Validity and Limits

- Eikonal/WKB: $|\nabla n|/n \ll k_0$. In strong field, integrate piecewise and avoid turning points.
- Paraxial expressions are used for closed-form deflection; otherwise evaluate $\alpha^\rightarrow = \int \nabla \perp \ln n$ ds numerically on $n(x)$.
- No new dimensional scales are introduced beyond $S_0 = \hbar$; c and G are acceptance locks.

7. Replication Examples

Solar grazing deflection: with $b \approx R_\odot$ and M_\odot ,

$$|\alpha| \approx 4 G M_\odot / (R_\odot c^2) \approx 1.75 \text{ arcsec}$$

Shapiro delay at superior conjunction: logarithmic microsecond-scale shift with the same η_0 .

Near-Earth gravitational redshift: $\Delta v/v \approx g h / c^2$ at laboratory heights (Pound–Rebka scale).

Appendix Pointers

- Eikonal working: gradients of $U = A e^{\{i k_0 S\}}$; transport; derivation of $d/ds(n t) = \nabla n$.
- Deflection integral along straight-line parameterization ($x=b, z=s$).
- Delay integral and endpoint geometry for the logarithm.
- Time-dilation mapping to $\sqrt{1 + 2\varphi}$.
- Periapsis advance via effective potential.
- Rotators: odd-parity term and far-field limit.

Bottom Line

One ontology, one scale ($S_0 = \hbar$), one cost map $n(x)$. The once-set factor that matches solar deflection also fixes delay and clock-shift. No extra knobs appear mid-derivation; kernels are short and SI-locked.

Appendix Setup (prep for detailed Math Appendix)

The companion Appendix carries the line-by-line evaluations alluded to above, using the same symbols and sign conventions stated here. Pointers:

- Rotators: odd-parity correction (far-field Lense–Thirring limit) and bench-style bounds.
- Periapsis precession: effective potential constructed from $d/ds(n t) = \nabla n$ and small-eccentricity expansion.
- Time-dilation: local cost-per-time ratio leading to $\sqrt{1 + 2\varphi}$ and $\Delta v/v \approx -\Delta\varphi$.
- Delay kernel: endpoint geometry for the standard log expression; relation to ψ in the thin-lens dictionary.
- Deflection integral in a $1/r$ potential: straight-line parameterization ($x = b, z = s$) and extraction of the logarithmic term.